

Today's outline - September 11, 2024



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- Refraction & reflection introduction

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- Boundary conditions at an interface

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Reading Assignment: Chapter 3.5–3.8

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- Refraction & reflection introduction
- Boundary conditions at an interface
- The Fresnel equations
- Reflectivity and Transmittivity
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Reading Assignment: Chapter 3.5–3.8

Homework Assignment #02:
Problems on September
due Monday, September 16, 2024

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- Refraction & reflection introduction
- Boundary conditions at an interface
- The Fresnel equations
- Reflectivity and Transmittivity
- Normalized q -coordinates

Reading Assignment: Chapter 3.5–3.8

Homework Assignment #02:
Problems on September
due Monday, September 16, 2024

Homework Assignment #03:
Chapter 3: 1,3,4,6,8
due Monday, September 30, 2024

Refraction & reflection of x-rays



X-rays can be treated like light when interaction with a medium. However, unlike visible light, the index of refraction of x-rays in matter is very close to unity:

Refraction & reflection of x-rays



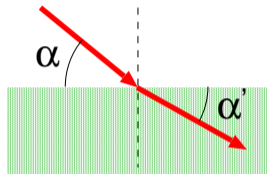
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$$n = 1 - \delta + i\beta, \quad \text{with } \delta \sim 10^{-5}$$

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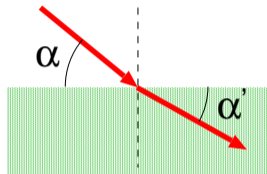


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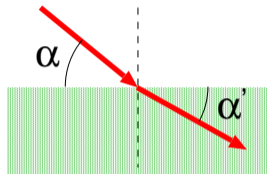
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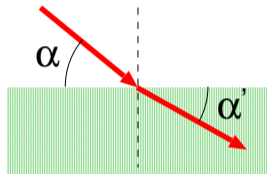
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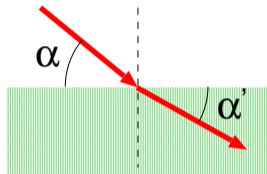
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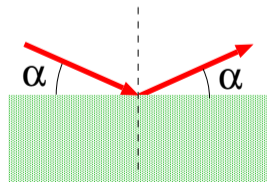


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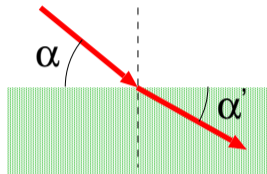
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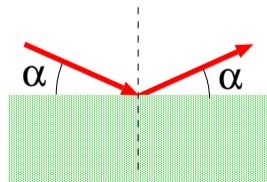


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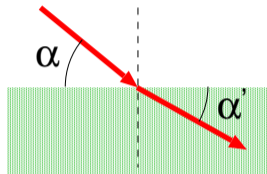


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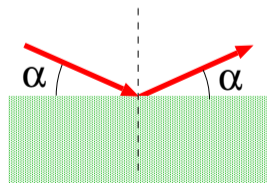


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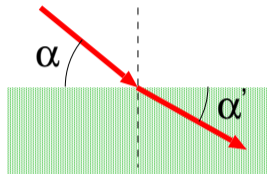
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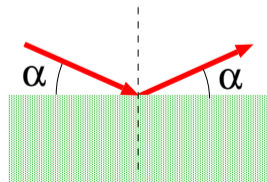


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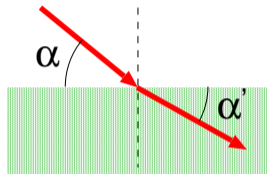
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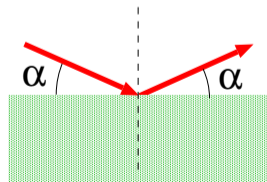


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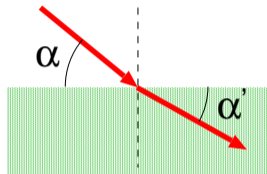
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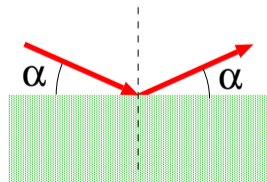


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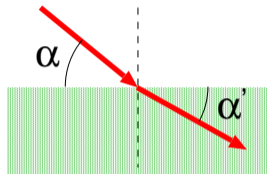
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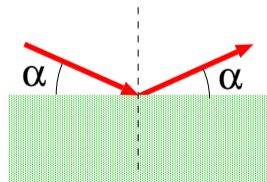


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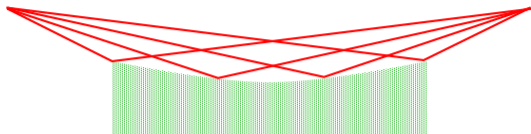


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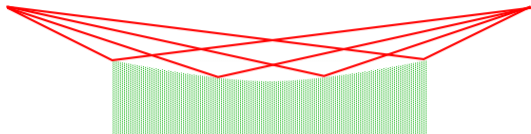
$$\delta = \frac{\alpha_c^2}{2} \quad \rightarrow \quad \alpha_c = \sqrt{2\delta}$$

Uses of total external reflection



X-ray mirrors

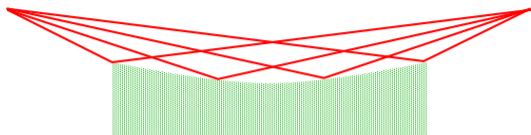
Uses of total external reflection



X-ray mirrors

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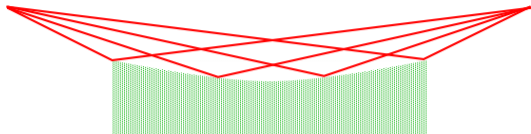
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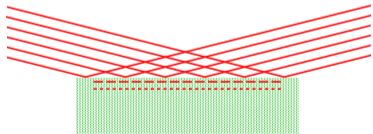
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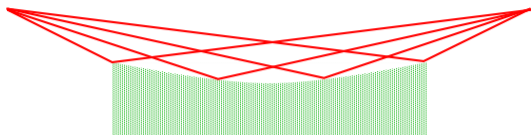
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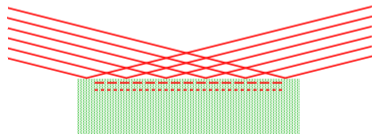
Evanscent wave experiments

Uses of total external reflection



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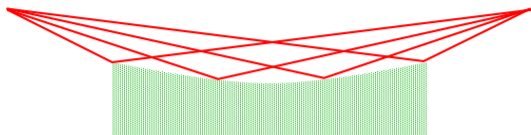
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Evanescent wave experiments

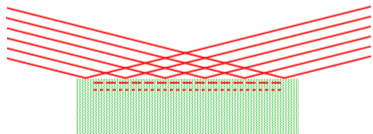
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Uses of total external reflection



X-ray mirrors

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Evanscent wave experiments

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Refractive index in the x-ray region



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Refractive index in the x-ray region



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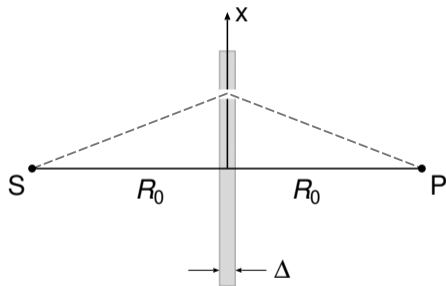
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Initially assume that all interfaces are perfectly flat and ignore all absorption processes.

Thin plate response - scattering approach



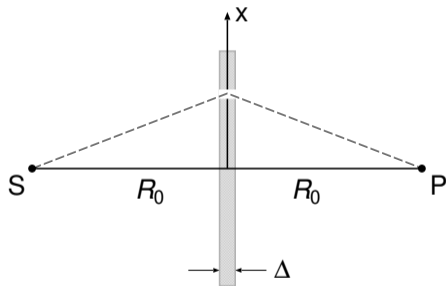
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Thin plate response - scattering approach

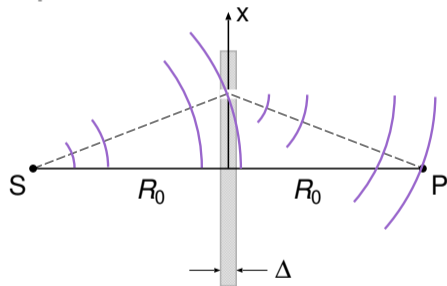
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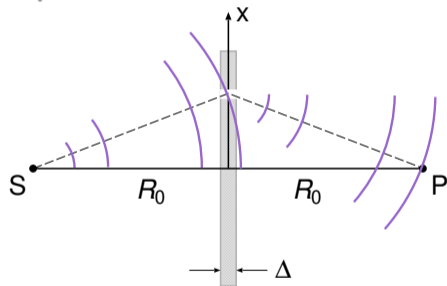
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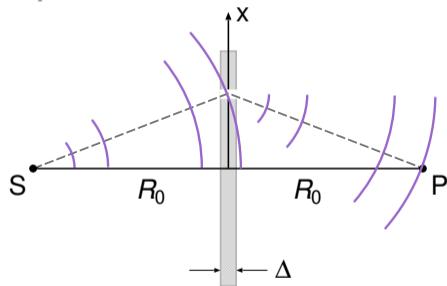


The plate has electron density ρ and the volume $\Delta dx dy$ contains $\rho \Delta dx dy$ electrons which scatter the x-rays.

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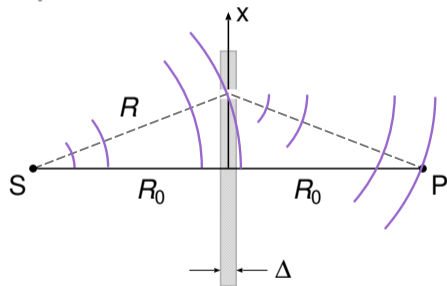


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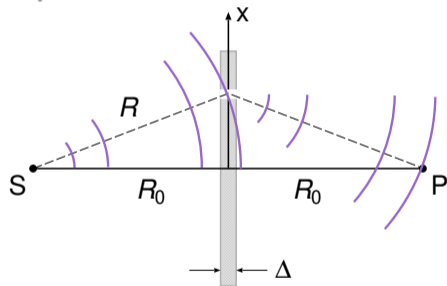
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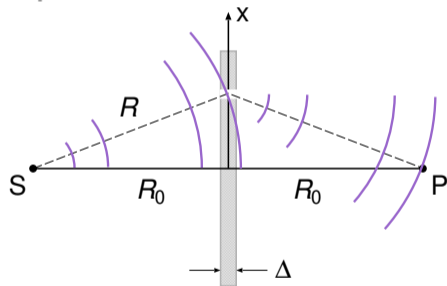
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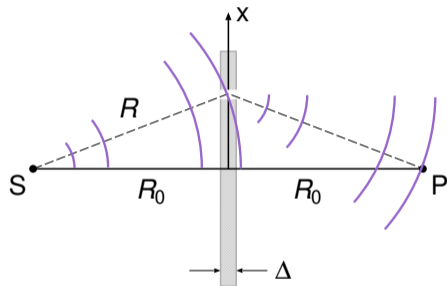
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Thin plate response - scattering approach



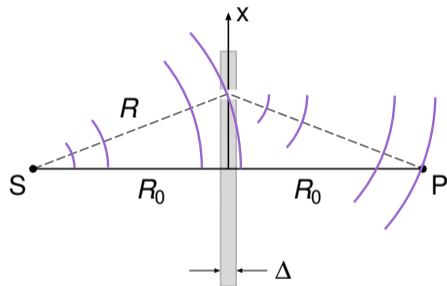
R is also the distance between the scattering volume and P so, a wave (x-ray) which travels from $S \rightarrow P$ through the scattering volume will have an extra phase shift



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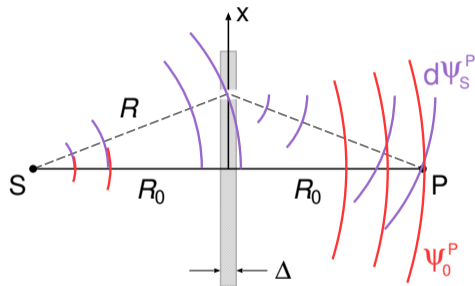


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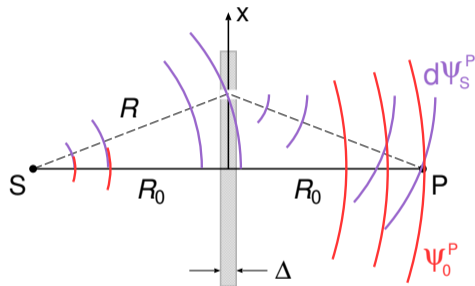
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compared to a wave which travels directly along the z-axis.

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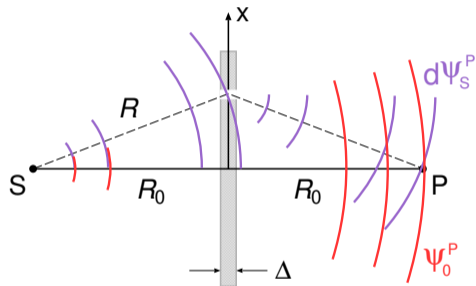
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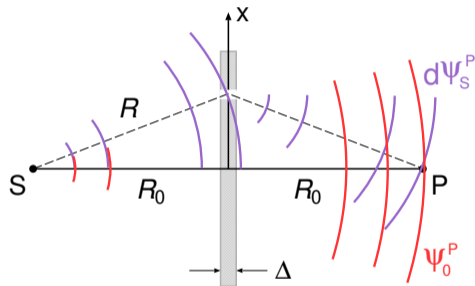
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$$d\psi_S^P \approx \left(\frac{e^{ikR_0}}{R_0} \right)$$



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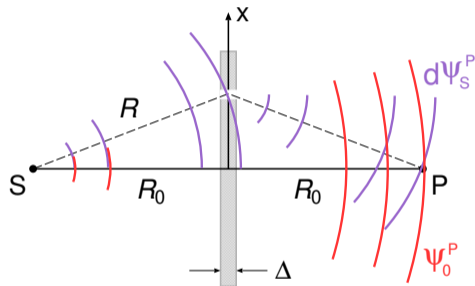
compared to a wave which travels directly along the z -axis. The wave which is scattered through the volume will have the form

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Thin plate response - scattering approach



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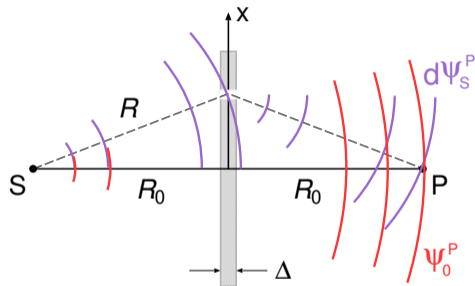
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Thin plate response - scattering approach



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Thin plate response - scattering approach



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Thin plate response - scattering approach



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Thin plate response - scattering approach



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Thin plate response - scattering approach



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Thin plate response - scattering approach



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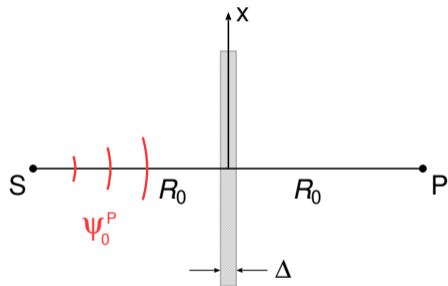
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Thin plate response - refraction approach

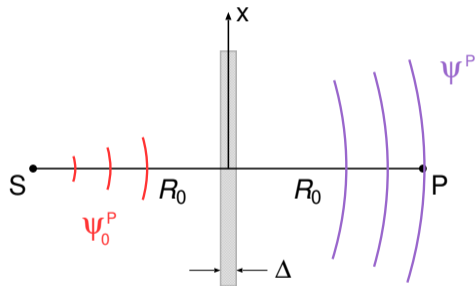
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Thin plate response - refraction approach



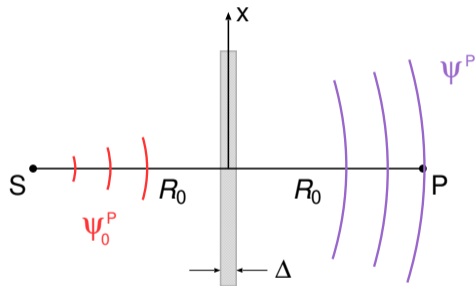
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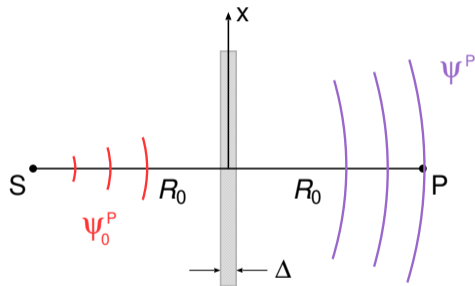


The phase shift depends on the thickness and the difference between the index of refraction of the medium and that of vacuum

Thin plate response - refraction approach



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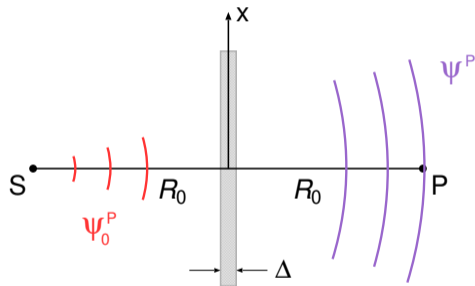
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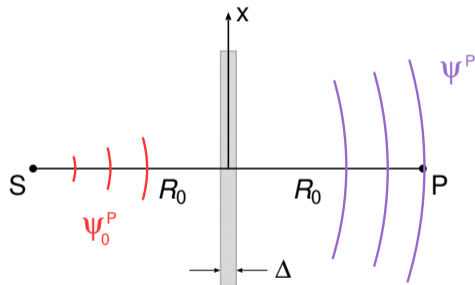
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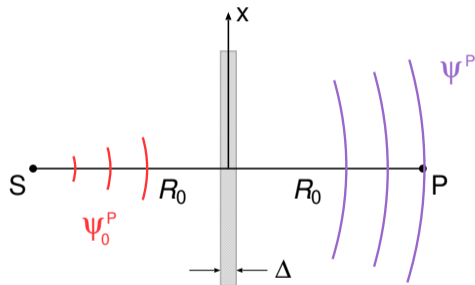
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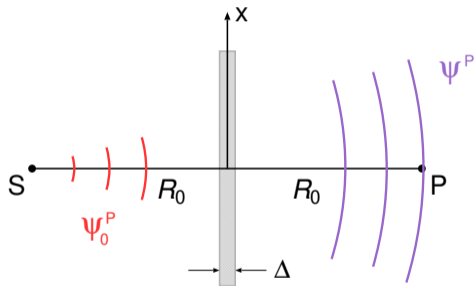
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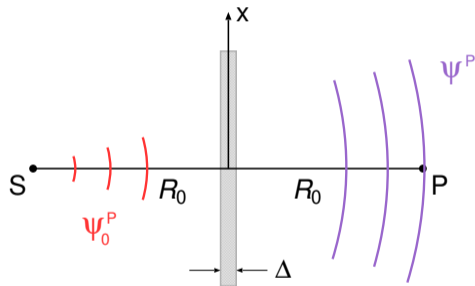
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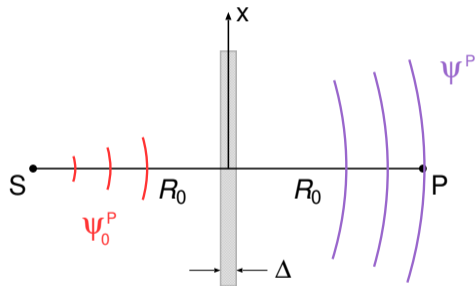
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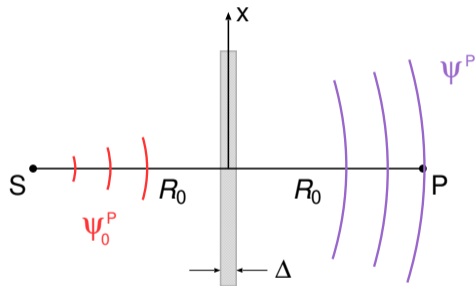
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Thin plate response - refraction approach



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Calculating n



We can now compare the expressions obtained by the scattering and refraction approaches.

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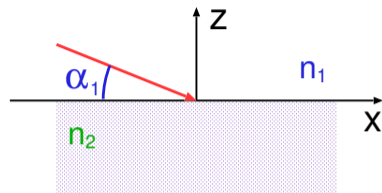
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Index of refraction & critical angle



Consider an x-ray incident on an interface at angle α_1 to the surface

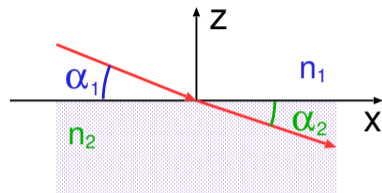


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Index of refraction & critical angle



Consider an x-ray incident on an interface at angle α_1 to the surface which is refracted into the medium of index n_2 at angle α_2 .



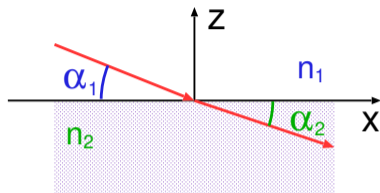
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Applying Snell's Law



$$n_2 \cos \alpha_2 = n_1 \cos \alpha_1$$

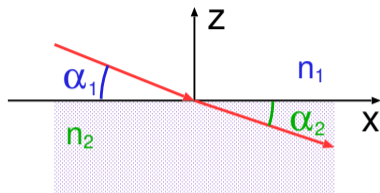
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Index of refraction & critical angle



Consider an x-ray incident on an interface at angle α_1 to the surface which is refracted into the medium of index n_2 at angle α_2 .

Applying Snell's Law, and assuming that the incident medium is "vacuum" ($n_1 = 1$).



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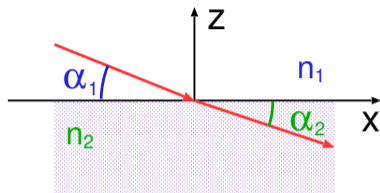
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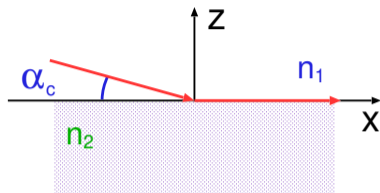
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When the incident angle becomes small enough, there will be total external reflection and $\cos \alpha_2 \equiv 1$

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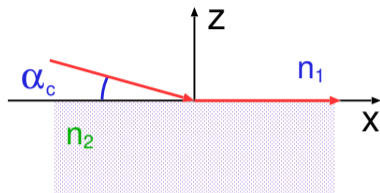
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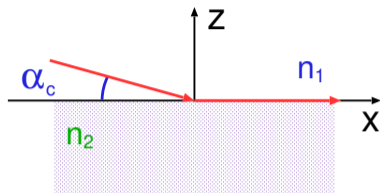
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Index of refraction & critical angle



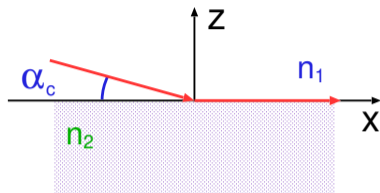
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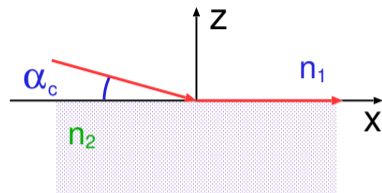
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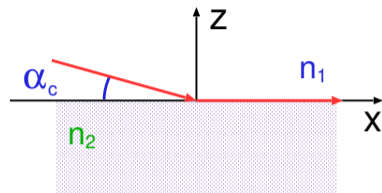


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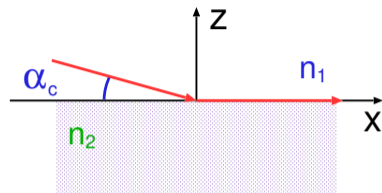


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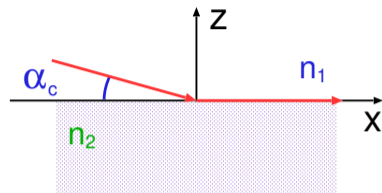


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Connection to atomic scattering



So far, we have made the assumption that the charge distribution is uniform. We know that this is not correct, and that usually electron charge distributions are those of the atoms making up the solid.

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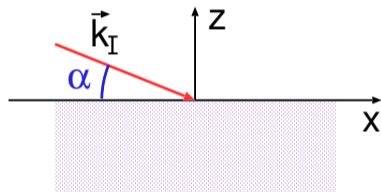
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Electromagnetic boundary conditions



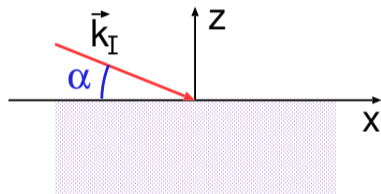
Maxwell's equations require that an electromagnetic wave and its derivative be continuous in all directions at any interface. This condition places restrictions on the waves which exist at any interface:



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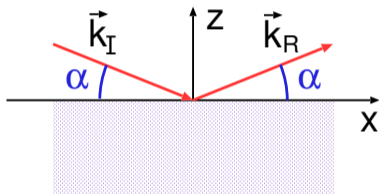


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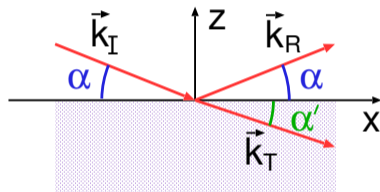
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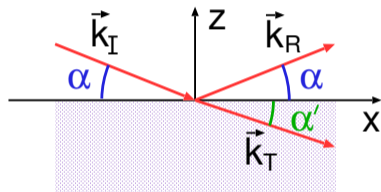
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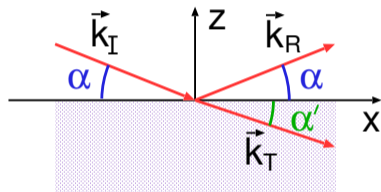
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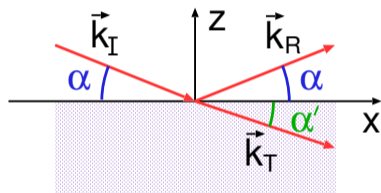
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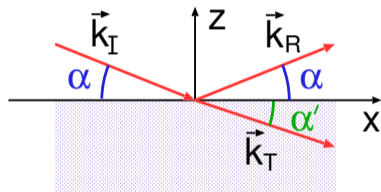
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Electromagnetic boundary conditions



Maxwell's equations require that an electromagnetic wave and its derivative be continuous in all directions at any interface. This condition places restrictions on the waves which exist at any interface:



which leads to conditions on the amplitudes and the wave vectors of the waves at $z = 0$. Taking vector components:

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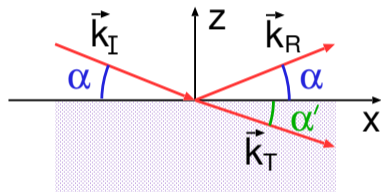
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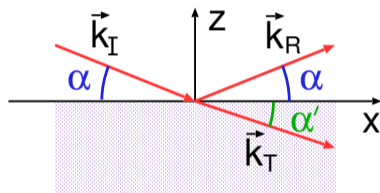
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Parallel projection & Snell's Law



Starting with the equation for the parallel projection of the field on the surface and noting that

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$$\begin{aligned} (a_I + a_R)n \cos \alpha' &= (a_I + a_R) \cos \alpha \\ \cos \alpha &= n \cos \alpha' \end{aligned}$$

Recalling that $\alpha_c = \sqrt{2\delta}$

$$1 - \frac{\alpha^2}{2} = (1 - \delta + i\beta) \left(1 - \frac{\alpha'^2}{2}\right) \quad \longrightarrow \quad \alpha^2 = \alpha'^2 + \alpha_c^2 - 2i\beta$$

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$$q = \frac{Q}{Q_c} \approx \frac{2k}{Q_c} \alpha \quad q' = \frac{Q'}{Q_c} \approx \frac{2k}{Q_c} \alpha'$$

q is a convenient parameter to use because it is a combination of two parameters which are often varied in experiments, the angle of incidence α and the wavenumber (energy) of the x-ray, k .

Defining equations in q



Start with the reduced version of Snell's Law

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Start with the reduced version of Snell's Law and multiply by a $1/\alpha_c^2 = (2k/Q_c)^2$.

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Defining equations in q

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$$r = \frac{q - q'}{q + q'}$$

$$t = \frac{2q}{q + q'}$$